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Key Points:

- SPE-enhanced atmospheric NO_y does not account for NO₃⁻ spikes in ice cores
- NO₃⁻ in ice cores is a poor proxy for historical SPEs

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Nitrate ion spikes in ice cores not suitable as proxies for solar proton events

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Abstract Nitrate ion spikes in polar ice cores are contentiously used to estimate the intensity, frequency, and probability of historical solar proton events, quantities that are needed to prepare for potentially society-crippling space weather events. We use the Whole Atmosphere Community Climate Model to calculate how large an event would have to be to produce enough odd nitrogen throughout the atmosphere to be discernible as nitrate peaks at the Earth's surface. These hypothetically large events are compared with probability of occurrence estimates derived from measured events, sunspot records, and cosmogenic radionuclides archives. We conclude that the fluence and spectrum of solar proton events necessary to produce odd nitrogen enhancements equivalent to the spikes of nitrate ions in Greenland ice cores are unlikely to have occurred throughout the Holocene, confirming that nitrate ions in ice cores are not suitable proxies for historical individual solar proton events.

1. Introduction

Direct measurements of solar energetic particles associated with solar flares and coronal mass ejections have only been available since the 1960s. The spectral shape and fluence of solar proton events (SPEs) prior to the launch of cosmic ray detectors aboard satellites can be inferred using records from ground-based ionization chambers and neutron monitors, available since the 1930s, or visual observations of sunspots and solar flares extending back to the 1600s. There has also been considerable effort to identify signatures of solar events within paleoclimate records in order to quantify the historical frequency and intensity of events, necessary to predict future solar activity and to study the connection between solar variability and climate.

The use of nitrate ion measurements in Antarctic and Greenland ice cores has been proposed as a method to estimate the intensity, frequency, and probability of individual solar proton events [e.g., *Zeller and Dreschhoff*, 1995; *McCracken et al.*, 2001; *Kepko et al.*, 2009], to forecast future solar activity [e.g., *Kim et al.*, 2009; *Barnard et al.*, 2011; *Riley*, 2012], and to determine upper limits for the fluence and energies of solar cosmic rays [e.g., *Shea et al.*, 2006; *Townsend et al.*, 2006; *Shea and Smart*, 2012; *Miroshnichenko and Nymmik*, 2014]. However, the use of nitrate as a proxy has proven controversial as a result of (1) the inability of models to reproduce enough atmospheric odd nitrogen or the transport mechanisms required to explain nitrate ion peaks in surface snow and ice [e.g., *Legrand et al.*, 1989; *Legrand and Kirchner*, 1990; *Duderstadt et al.*, 2014], (2) chemical correlations within ice cores pointing to alternative tropospheric sources [e.g., *Legrand and Mayewski*, 1997; *Wolff et al.*, 2008, 2012], and (3) the recognition of postdepositional processes involving nitrate ions during sequestration within ice cores [e.g., *Dibb and Whitlow*, 1996; *Dibb and Jaffrezo*, 1997; *Röthlisberger et al.*, 2002]. For these reasons, *Schrijver et al.* [2012] exclude nitrate as a proxy for SPEs in their analysis of the frequency of extreme solar events. Nonetheless, the use of nitrate as a proxy for SPEs persists, exemplified by the *Smart et al.* [2014] counterargument that only very high resolution nitrate ion measurements can capture SPEs, such as the analysis of the Greenland Ice Sheet Project (GISP2-H) and Boston University (BU) ice cores from Summit, Greenland.

Solar protons enhance odd nitrogen ($\text{NO}_y = \text{N} + \text{NO}_2 + \text{NO}_3 + 2\text{N}_2\text{O}_5 + \text{HNO}_3 + \text{HO}_2\text{NO}_2 + \text{ClONO}_2 + \text{BrONO}_2$) and odd hydrogen ($\text{HO}_x = \text{H} + \text{OH} + \text{HO}_2$) in the polar upper stratosphere and mesosphere through the ionization and dissociation of molecular N₂ and O₂ [e.g., *Crutzen et al.*, 1975; *Jackman et al.*, 1980; *Solomon et al.*, 1981]. Strong downward transport within the winter polar vortex brings enhanced NO_y to the lower

stratosphere; the NO_y is potentially deposited at the surface as nitrate ions through the gravitational settling of polar stratospheric cloud particles or through stratosphere-troposphere exchange followed by wet or dry deposition. Satellite observations of NO_2 have confirmed the SPE enhancement of NO_y and its subsequent descent [e.g., López-Puertas *et al.*, 2005; Randall *et al.*, 2005; Jackman *et al.*, 2009; Randall *et al.*, 2009]. However, modeling studies have been unable to produce sufficient enhancement of odd nitrogen from SPEs in the upper atmosphere to account for spikes seen in surface snow and ice [e.g., Jackman *et al.*, 1990; Vitt *et al.*, 2000; Calisto *et al.*, 2012; Duderstadt *et al.*, 2014]. In addition, given its long stratospheric residence time, diabatic recirculation, and dilution to lower latitudes, SPE-enhanced NO_y in the mesosphere and stratosphere from individual events is unlikely to be discernible at the surface [e.g., Legrand *et al.*, 1989; Legrand and Kirchner, 1990]. Most of these modeling studies consider energies limited to 1–300 MeV, corresponding to proton channels measured by the energetic particle sensor (EPS) instruments on GOES satellites [e.g., Jackman *et al.*, 1980, 2005, 2008]. There have been several efforts to model the atmospheric impact of high-energy solar protons by also including satellite high-energy proton and alpha detector (HEPAD) instrument observations (300 to >700 MeV), inferences from ground level enhancements (GLEs), and extrapolations of satellite observations to higher energies with functional fits [e.g., Rodger *et al.*, 2008; Calisto *et al.*, 2012, 2013; Kovaltsov and Usoskin, 2014]. These studies focus on the effect of large SPEs on stratospheric ozone, temperature, and winds. The question remains whether higher-energy SPEs, evidenced by GLEs from neutron monitors and muon detectors, produce enough NO_y directly in the troposphere and lower stratosphere to account for spikes of nitrate ions at the surface.

Our work uses the Whole Atmosphere Community Climate Model (WACCM) to consider higher proton energies with the goal of estimating the fluence and spectra of solar protons that would be necessary to produce discernible nitrate ion spikes at the surface. We place our results in the context of the probability of occurrence of large SPEs derived from spacecraft, GLEs, and cosmogenic radionuclide archives. How large do SPEs have to be in order to produce enhancements of NO_y equivalent to the nitrate spikes attributed to SPEs in the GISP2-H and BU ice cores? Do these SPEs occur frequently enough for nitrate ion records to provide meaningful statistics of historical SPEs?

Section 2 of this manuscript discusses measurements of solar protons and strategies for extrapolating proton flux to higher energies. Section 3 describes the Whole Atmosphere Community Climate Model (WACCM) and model scenarios. Section 4 presents model results in comparison with nitrate ion spikes in snow and ice. Section 5 discusses results in the context of the historical probability of solar extreme events. Section 6 summarizes conclusions regarding the efficacy of using nitrate peaks in ice cores to study historical SPEs.

2. The Fluence and Spectra of Solar Proton Events

A complete energy spectrum for SPEs can be constructed by extrapolating satellite proton flux spectra to relativistic energies, using ground level neutron or muon measurements to infer solar proton flux, or a combination of both. For this work, we rely on highly documented SPEs as a basis for constructing hypothetically large events. We scale these SPEs in fluence and duration to produce atmospheric odd nitrogen enhancements of relative magnitudes similar to nitrate ion peaks attributed to SPEs in ice cores.

GOES energetic particle sensors (EPSs) measure seven logarithmically spaced proton differential energy bands ranging from 1 to ~500 MeV [NASA, 1996]. High-energy proton and alpha particle detectors (HEPADs) measure bands from 350 to 700 MeV, with an additional integral energy channel above 700 MeV. Based on the discussions of Smart and Shea [1999] and NASA [2006], we do not include channel P7 of EPS in our spectrum fits because the P7 spectral width (165–500 MeV) results in significant uncertainty regarding the choice of a midpoint energy. For the HEPAD channel P11, measuring the integral flux of energies greater than 700 MeV, we assume a midpoint energy of 1000 MeV following recommendations by Smart and Shea [1999]. While Smart and Shea [1999] also outline a strategy for compensating for side and rear penetration of high-energy particles to lower energy channels, we have chosen not to include these corrections as they appear to be minimal and limited to specific events [Mottl and Nymmik, 2007] and are not consistent with our efforts to estimate the maximal potential enhancement of odd nitrogen from SPEs.

Optimization methods use ground level measurements from neutron monitors to infer solar energetic particle spectra at the top of the atmosphere. We rely on the proton flux from these studies to model the 23 February

1956 SPE, an event widely used to estimate the strength of the 1859 Carrington white flare event as well as to predict future “worst-case scenarios” [e.g., *Feynman et al.*, 1990; *Shea and Smart*, 1990; *Webber et al.*, 2007]. Similar methods are also used to identify functions for extrapolating satellite observations to higher energies.

Past studies have used a variety of spectral shapes to construct proton flux energy distributions from measurements. Prior WACCM experiments use exponential fits to interpolate between proton flux observations ranging from 1 to 300 MeV [*Jackman et al.*, 2008, 2009, 2011; *Funke et al.*, 2011; *Duderstadt et al.*, 2014]. Other studies rely on Bessel functions [*McGuire and von Rosenvinge*, 1984], Weibull functions [e.g., *Xapsos et al.*, 2000; *Kim et al.*, 2009] and power laws with exponential rollover [e.g., *Ellison and Ramaty*, 1985; *Mazur et al.*, 1992; *Mewaldt et al.*, 2005]. The double power law function proposed by *Band et al.* [1993] appears to provide the best fit to events during the satellite era [e.g., *Mewaldt et al.*, 2005, 2012; *Tylka et al.*, 2006; *Atwell et al.*, 2011] and is particularly well suited to energies above 100 MeV [e.g., *Mottl et al.*, 2001; *Mewaldt et al.*, 2012]. This Band function results in distinct power law parameters for low and high energies with a break (aka “knee,” “bend point,” or “spectrum turnover”) at energies close to 30 MeV [e.g., *Mewaldt et al.*, 2005; *Wang and Wang*, 2006; *Mewaldt et al.*, 2012].

We interpolate observed SPE spectra to 60 logarithmically based energy levels between 1 and 700 MeV, using an exponential fit between each measured EPS and HEPAD energy level. We then extrapolate to energies beyond 700 MeV by fitting a power law to channels above P5 (40–80 MeV), beyond the break energies typically seen in double power law spectra [*Mewaldt et al.*, 2012]. Questions remain concerning the shape of the proton spectra at relativistic energies [e.g., *Gopalswamy et al.*, 2012; *Mewaldt et al.*, 2012; *Miroshnichenko et al.*, 2013]. For example, *Mewaldt et al.* [2012] identify a second spectral rollover around 500 MeV based on GLE data. *Miroshnichenko and Nymmik* [2014] also point out overestimates to the high-energy tails of early reconstructions from GLEs such as the 23 February 1956 event. Direct measurements of protons with rigidities ranging from 1 GV to 1.2 TV by the Payload for Antimatter Matter Exploration and Light-nuclei Astrophysics mission [*Adriani et al.*, 2011] and 1 GV to 1.8 TV by the Alpha Magnetic Spectrometer (AMS) on the International Space Station [*Aguilar et al.*, 2015] will allow more accurate estimates of higher energies. The spectral indices for the high-energy tails of hard spectrum hypothetical SPEs in our WACCM simulations are consistent with preliminary AMS data from 2011 to 2013 [*Aguilar et al.*, 2015]. On the basis of current knowledge, our high-energy tail extrapolations are consequently considered an upper bound, emphasizing the goal of estimating the maximum amount of odd nitrogen that could be produced in the lower atmosphere during an SPE.

The SPE proton flux reaching the top of the atmosphere evolves with time [e.g., *Wang*, 2009; *Vashenyuk et al.*, 2011; *Miroshnichenko et al.*, 2013], usually beginning with a “prompt” component that is impulsive, has anisotropic pitch angles, and exhibits a “hard” (high-energy) exponential energy spectrum. A subsequent “delayed” component has a smoother time profile, is close to isotropic, and has a “soft” (low-energy) power law spectrum. The majority of high-energy particles are found during the first 20 h of the event [*Li et al.*, 2003]. In this work, we capture this time evolution in spectral shape by basing our ion pair calculations on 30 min average differential energy flux observations, in contrast to the daily averages used in prior WACCM simulations.

Our construction of hypothetical extreme SPEs from known SPEs assumes a uniform scaling across all energies and preserves the time evolution of the observed spectral shape. This strategy is based on the study of *Crosby* [2009], arguing that solar extreme events can be viewed as part of a global distribution originating from similar solar and heliospheric mechanisms as opposed to outliers with peculiar characteristics. However, in the absence of direct measurements of the characteristics of SPEs at the high fluences discussed in this paper, the spectral shape and evolution of our hypothetical constructions remain uncertain.

3. WACCM Model

The Whole Atmosphere Community Climate Model (WACCM) is a three-dimensional Earth system model specifically designed to study the impact of upper atmospheric processes on global climate. WACCM is part of the Community Earth System Model (CESM) at the National Center for Atmospheric Research (NCAR) and combines the Community Atmosphere Model version 5, the thermosphere-ionosphere-mesosphere electrodynamics general circulation model, and the Model for Ozone and Related chemical Tracers [*Garcia et al.*, 2007; *Kinnison et al.*, 2007; *Emmons et al.*, 2010; *Marsh et al.*, 2013; *Neale et al.*, 2013]. By resolving the entire

stratosphere and mesosphere, WACCM is ideally suited for modeling SPE production of HO_x and NO_x in the upper atmosphere [e.g., *Jackman et al.*, 2008, 2009, 2011; *Funke et al.*, 2011; *Duderstadt et al.*, 2014].

For primary particles with energies of 1–300 MeV, ion pair production rates as a function of pressure are calculated using the parameterization of *Jackman et al.* [1980, 2005, 2008] based on the range data of *Sternheimer* [1959]. Higher energies use ionization yield functions from *Usoskin et al.* [2010] calculated using the CRAC:CRIL (Cosmic Ray-induced Atmospheric Cascade: Application for Cosmic Ray Induced Ionization) particle transport model. The Monte Carlo methods of the CRAC:CRIL reaction cascade resolve species and processes below 1 g cm⁻² (~50 km) [*Usoskin et al.*, 2004; *Usoskin and Kovaltsov*, 2006; *Usoskin et al.*, 2010] and account for the effects of nuclear processes and secondary particles.

Jackman et al. [1980, 2005] describe WACCM parameterizations for HO_x and NO_x production. WACCM relies on a look-up table to calculate the production of HO_x, resulting from the ionization of N₂ and O₂ followed by water cluster ion formation [*Solomon et al.*, 1981]. Regarding NO_x production in WACCM, the ionization and dissociation of N₂ and O₂ produce 1.25 N per ion pair [*Porter et al.*, 1976], subsequently partitioned into 45% ground state N(⁴S) and 55% excited state N(²D), with the excited state determining NO production [*Rusch et al.*, 1981]. These simulations assume uniform HO_x and NO_x production at geomagnetic latitudes greater than 60°.

Our WACCM experiments have a horizontal resolution of 1.9° latitude by 2.5° longitude with 88 vertical layers extending to ~140 km. These simulations use the Specified Dynamics (SD) version of WACCM. Meteorological fields in the model from 0 to 40 km are nudged to data from the NASA Global Modeling and Assimilation Office (GMAO) Modern-Era Retrospective Analysis for Research and Application meteorological reanalysis [*Rienecker et al.*, 2011]. There is a linear transition in the nudging from 40 to 50 km, above which the model is free running. To study the sensitivity of nitrate production to solar proton spectral shape and fluence, we reproduce several known SPEs assuming a start date of 18 December 2004. The Arctic winter of 2004–2005 was chosen because it was unusually cold, with strong diabatic descent and limited mixing to lower latitudes during December and January [*Manney et al.*, 2006]. Any stratospheric NO_y produced by an SPE under these meteorological conditions would remain well sequestered within the polar vortex with relatively unmixed descent for at least 6 weeks. SD-WACCM simulations for the 2004–2005 polar winter also show good correlation with satellite measurements of ozone and related constituents [*Brakebusch et al.*, 2013]. While *Brakebusch et al.* [2013] conclude that SD-WACCM overestimates mixing across the vortex edge in the latter half of the season, the model captures vertical transport well within the low-middle stratosphere, especially during December and January. This study does not address the sensitivity to dynamical conditions and is limited to the Northern Hemisphere. Although the 2004–2005 Arctic winter was cold, it is likely that conditions even more conducive to descent occurred in the past. However, even if all the NO_x produced by the hypothetical SPEs was sequestered in the polar vortex, the conclusions drawn here would not change.

WACCM includes the simulation of polar stratospheric clouds (PSCs). Although not as significant in the Arctic as the Antarctic, PSCs play a role in atmospheric chemistry during the 2004–2005 winter as a result of the cold, stable vortex early in the winter [*Brakebusch et al.*, 2013]. WACCM addresses heterogeneous reactions on sulfate aerosols, nitric acid trihydrate (NAT), supercooled ternary solution, and water ice associated with PSCs [*Kinnison et al.*, 2007]. Including simulations of PSCs is important for studying nitrate deposition due to the uptake of nitrate on PSCs followed by gravitational settling of PSCs to the troposphere. However, it is important to keep in mind that PSCs are infrequent in the Arctic and occur only inside the polar vortex during cold winters, with average occurrence frequencies of less than 6% [*Fromm et al.*, 2003].

4. Results

4.1. WACCM Experiment Setup

Prior global model simulations of SPEs have failed to produce enough NO_y in the upper atmosphere to account for nitrate ion spikes observed in snow and ice, generally showing enhancements of only 10–20% in the lower stratosphere during the months following large SPEs [e.g., *Vitt et al.*, 2000; *Jackman et al.*, 2008, 2009; *Duderstadt et al.*, 2014]. *Duderstadt et al.* [2014] conducted WACCM simulations for the 9 November 2000 SPE, the sixth largest event in the past 50 years in terms of NO_y production. The authors chose this event because of the availability of concurrent, year-round, daily snow surface measurements at

Table 1. Comparison of Solar Proton Events From 1956 to 2012

Solar Proton Event Date	NO _y Production ^a (Gmol)	GLE Number ^b	GLE Peak ^c (%)	F ₃₀ ^d (cm ⁻²)	F ₂₀₀ ^e (cm ⁻²)
23 Feb 1956	–	5	5000	1.8E + 09	1.2E + 08
10–17 Jul 1959	–	7	18	2.3E + 09	1.6E + 07
12–20 Nov 1960	–	10,11,12	150,100,–	9.0E + 09	1.0E + 08
2–7 Sep 1966	2.0	–	–	–	–
28 Jan to 1 Feb 1967	1.6	16	30	–	4.3E + 06
2–10 Aug 1972	6.0	24,25	79,17	5.0E + 09	1.7E + 07
13–26 Aug 1989	3.0	41	24	1.5E + 09	5.1E + 06
29 Sep to 3 Oct 1989	1.7	42	395	1.8E + 08	3.1E + 07
19–27 Oct 1989	11	43,44,45	90,190,95	4.3E + 09	9.3E + 07
21–28 May 1990	–	47,48,49,50	24,50,–,–	–	6.4E + 06
11–15 Jun 1991	–	51,52	12,57	–	8.3E + 06
14–16 Jul 2000	5.8	59	80	4.3E + 09	3.4E + 07
9–11 Nov 2000	4.3	–	–	3.1E + 09	–
15–18 April 2001	–	60,61	230,26	1.5E + 08	9.3E + 06
24–30 Sep 2001	3.3	–	–	1.2E + 09	–
5–7 Nov 2001	5.3	62	–	3.4E + 09	1.2E + 07
23–25 Nov 2001	2.8	–	–	8.5E + 08	–
28–31 Oct 2003	5.6	65,66	45,35	3.3E + 09	2.3E + 07
15–23 Jan 2005	1.8	68,69	–,5500	1.0E + 09	2.4E + 07
23–30 Jan 2012	1.9	–	–	5.5E + 08	–
7–11 Mar 2012	2.1	–	–	9.6E + 08	–

^aComputed NO_y production in middle atmosphere (gigamoles) [Jackman *et al.*, 2008, 2009, 2014].

^bGLE number is associated with the number assigned to each GLE recorded by neutron monitors since 28 February 1942.

^cMaximum peak percent increase for corresponding GLEs from Neutron Monitor Database as describe by McCracken *et al.* [2012]. The early data are mostly 5 min averages, with 1 min averages more frequent after 1990. Values less than 10% are represented by “–.”

^dFluence >30 MeV [Shea *et al.*, 2006; Smart *et al.*, 2006b; Reedy, 2006; Webber *et al.*, 2007; Reedy, 2015]. Values less than 1×10^9 cm⁻² or unknown to the authors are represented by “–.”

^eFluence >200 MeV [Kovaltsov *et al.*, 2014]. Values less than 1×10^6 cm⁻² or unknown to the authors are represented by “–.”

Summit, Greenland, including a comprehensive suite of ion measurements capable of distinguishing probable tropospheric sources of nitrate from possible SPE influences. While WACCM calculates a thin ~5 km layer of 5–10 ppbv SPE-enhanced NO_y at 25–30 km in the months following the 9 November 2000 event, the thick background pool of 10–15 ppbv NO_y in the lower stratosphere results in a total column atmospheric enhancement of less than 5%, not enough to explain the twofold to fivefold spikes in nitrate observed in surface snow. Justification for the use of NO_y column densities and polar vortex averages is described in detail in Duderstadt *et al.* [2014]. The 9 November 2000 event, however, was not the highest fluence event in past decades nor did it produce GLEs at the surface. In our current study, we expand upon this prior research by considering SPEs with larger fluences and higher-energy (“harder”) spectra, noting that daily snow measurement comparisons are not available for these events.

Table 1 compares several of the largest SPEs occurring since the 1950s (derived from Reedy [2006], Shea *et al.* [2006], Webber *et al.* [2007], Jackman *et al.* [2008, 2009, 2014], McCracken *et al.* [2012], and Kovaltsov *et al.* [2014]). The characteristics of events prior to the satellite era are inferred from ground level measurements of secondary particles [e.g., Feynman *et al.*, 1990; Shea and Smart, 1990]. Table 2 summarizes the modeling experiments presented in this paper. Our first set of WACCM experiments (section 4.2) compares several satellite era SPEs that have large fluences of protons with energy greater than 30 MeV (defined to be F₃₀), including the series of SPEs in October 1989, the 14 July 2000 “Bastille Day” event, the 9–11 November 2000 SPE, the 28–31 October 2003 “Halloween Storms,” and the 23 July 2012 SPE that missed Earth but was observed by the STEREO-A spacecraft. A second set of modeling experiments (section 4.3) evaluates odd nitrogen production by hard spectrum SPEs, specifically the events of 29 September 1989 and 20 January 2005 that resulted in the largest GLEs in the satellite era. Finally, we present simulations of the 4 August 1972 and 23 February 1956 SPEs (section 4.4), comparing WACCM calculations with other studies striving to use “Carrington-like” conditions to study worst-case scenarios. We place our WACCM results in

Table 2. Simulation Runs

Observed	Hypothetical High Fluence	Hypothetical Hard Spectra	Carrington Like
29 Sep 1989	19–27 Oct 1989 $\times 10$	29 Sep 1989 $\times 100$	23 Feb 1956 (scaled to $1.9 \times 10^{10} \text{ cm}^{-2} F_{30}$)
19–27 Oct 1989	19–27 Oct 1989 $\times 50$	29 Sep 1989 $\times 1000$	
14 Jul 2000	19–27 Oct 1989 $\times 100$	29 Sep 1989 10d $\times 100$	
9 Nov 2000	19–27 Oct 1989 10d $\times 10$	20 Jan 2005 $\times 100$	4 Aug 1972 (scaled to $1.9 \times 10^{10} \text{ cm}^{-2} F_{30}$)
26–31 Oct 2003		20 Jan 2005 $\times 1000$	
15–23 Jan 2005		20 Jan 2005 10d $\times 100$	
23 Jul 2012			

the context of nitrate ion spikes measured in the GISP2-H and BU ice cores from Summit, Greenland. For all simulations, we increase the flux and duration of events in order to determine the fluence and spectrum of an SPE necessary to produce enough odd nitrogen in the atmosphere to potentially be discernible as nitrate spikes at the surface.

Figures 1 and 2 show the differential energy spectra and ion pair production rates for several SPEs in this study. The vertical axis of Figure 1 represents the midpoint energy levels of the GOES proton channels (unequally spaced), and the black dashed lines denote 30 MeV and 200 MeV cutoffs corresponding to Table 1. The proton flux for the 23 July 2012 event uses hourly STEREO-A observations from 1 to 100 MeV extrapolated to 500 MeV based on the Predictions of Radiation from REleASE, EMMREM, and Data Incorporating the CRaTER, COSTEP, and other SEP measurements model [Schwadron, 2012; Joyce et al., 2013].

The 14 July 2000 event clearly has the largest F_{30} fluence during peak intensity, resulting in high rates of ionization from 40 to 70 km. The extended duration of the October 1989 and October 2003 events explains their large total F_{30} fluences as well, resulting in significant ionization throughout the upper stratosphere and mesosphere. The majority of ion pair production occurs above 20 km. Ionization below 20 km is evident during the 20 January 2005 SPE, and there are also high ionization rates at relatively low stratospheric altitudes during the 29 September 1989, 20 October 1989, and 14 July 2000 SPEs. These figures confirm the conclusion by Usoskin et al. [2011] and Mironova and Usoskin [2013] that ion pair production rates in the lower stratosphere are 1–2 orders of magnitude less than in the midstratosphere, even during hard spectrum events.

4.2. WACCM Simulations of High Fluence Events

Figure 3 presents NO_y enhancements from several high fluence events (F_{30}) averaged over the Arctic polar vortex throughout winter. All events are placed beginning on 18 December 2004. The boundary of the meandering polar vortex is objectively determined by identifying grid points within the stratosphere where scaled potential vorticity (SPV) on isentropic surfaces exceeds $1.4 \times 10^{-4} \text{ s}^{-1}$ [Dunkerton and Delisi, 1986], as described in further detail in Duderstadt et al. [2014]. This figure illustrates NO_y enhancements in the upper stratosphere and mesosphere during the events as well as diabatic descent within the cooling polar vortex throughout winter. This figure also shows background NO_y (without SPEs) in units of mixing ratio and number density, highlighting the pool of NO_y in the lower stratosphere formed from N_2O oxidation. This background concentration of NO_y can also be brought to the surface following tropopause folds (stratosphere-troposphere exchange) and denitrification on polar stratospheric cloud particles, suggesting that SPE enhancements of odd nitrogen would need to be much larger than this reservoir to be detected at the surface.

A comparison of total column density for simulations with and without SPEs provides a simple means of assessing SPE enhancements of odd nitrogen with respect to background. Figure 4 presents vortex-averaged total NO_y column densities for these large fluence events. All SPE NO_y enhancements remain below 15% relative to background levels, with localized enhancements (individual grid points within the polar vortex) below 20% (not shown).

We note that values of NO_y presented in this paper include both gas and condensed phase HNO_3 , allowing the study of the maximal potential for nitrate deposition from both gas phase reactions and heterogeneous reactions on PSCs as well as the effects of the formation, settling, and evaporation of PSCs. Including the impact of PSCs on SPE-enhanced NO_y is particularly important during the 2004–2005 winter, with WACCM simulations suggesting NAT PSC formation beginning in early December from 20 to 30 km, gradually extending to 15–30 km by the end

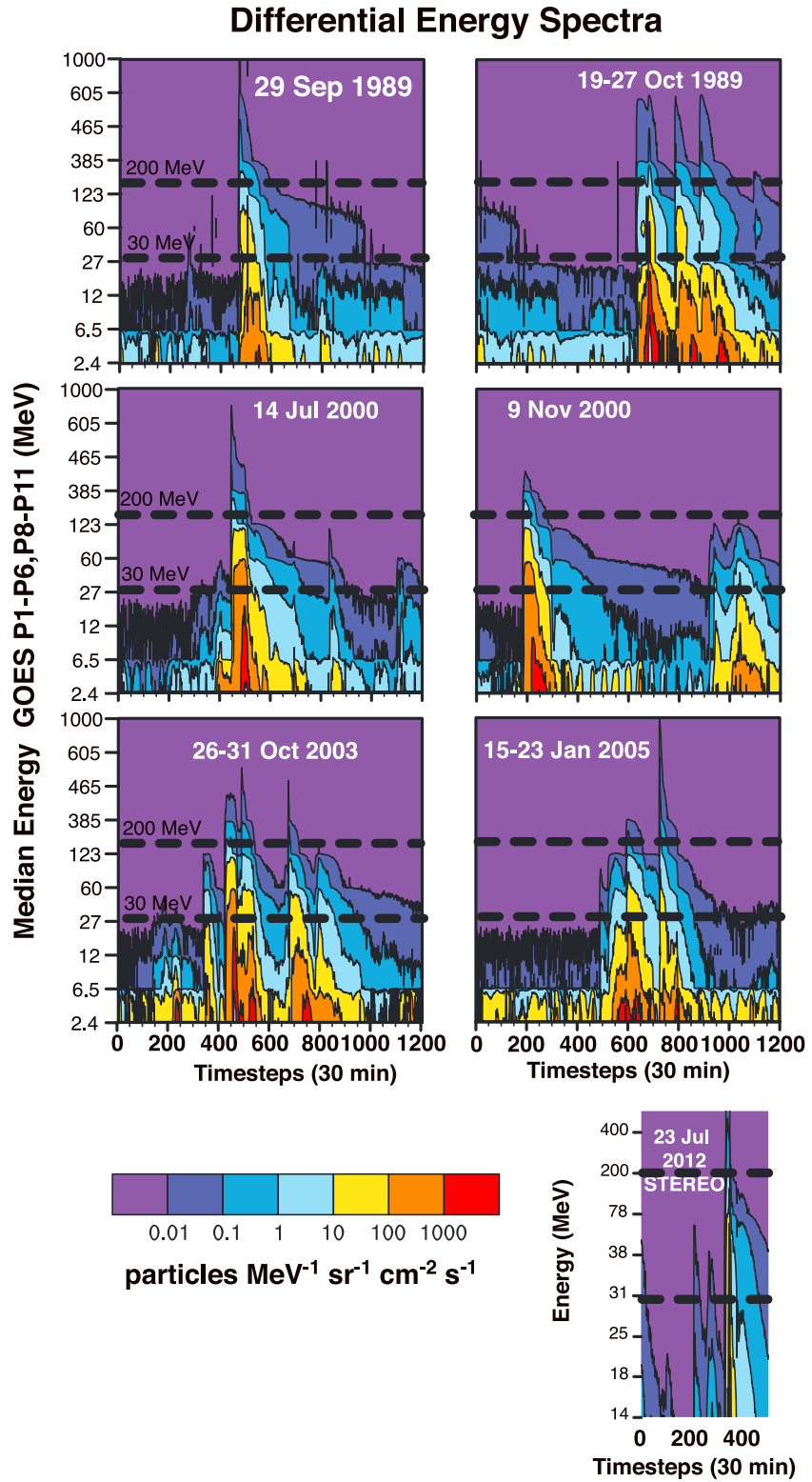


Figure 1. Differential energy flux as a function of GOES and STEREO energy level and time. Note that vertical axis refers to energy channel and is neither linear nor logarithmic. Dashed lines denote 30 MeV and 200 MeV energy levels corresponding to Table 1.

Ion Pair Production Rates

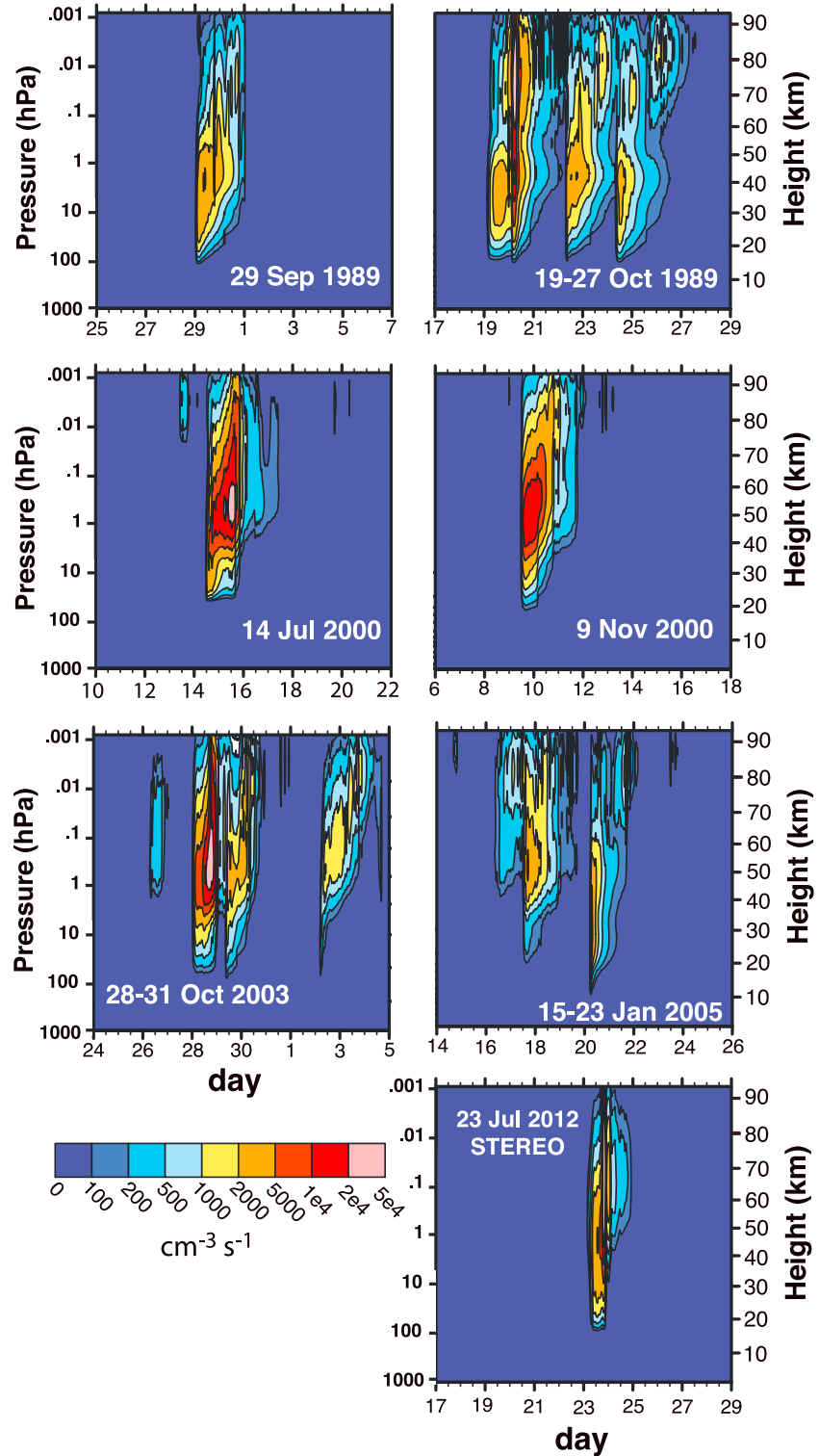


Figure 2. Ion pair production rates ($\text{cm}^{-3} \text{s}^{-1}$) as a function of altitude (pressure and height) and time (day of the month).

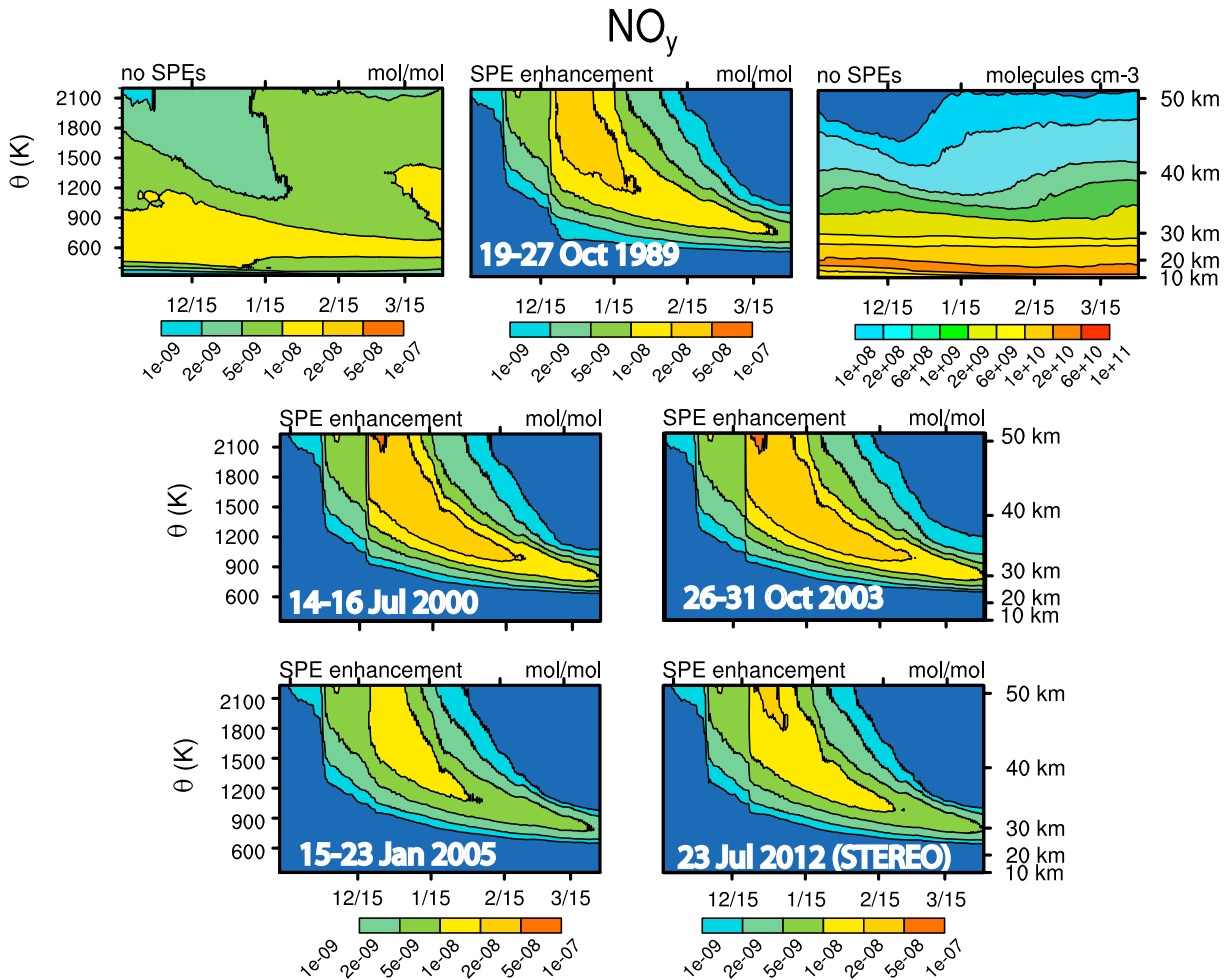


Figure 3. Time evolution of the Arctic vortex-averaged enhancement of NO_y (mole ratios) for SPEs placed in the winter of 2004–2005. (top left) Background NO_y (mole ratios) with no SPEs. (top right) Background NO_y (number density) with no SPEs. Dates are given as month/day.

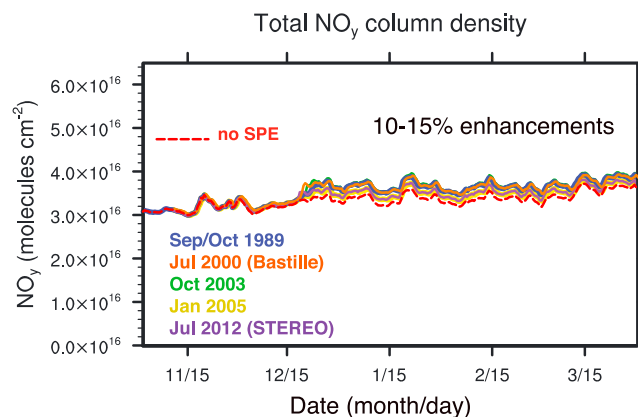


Figure 4. Arctic vortex-averaged total NO_y column density (molecules cm^{-2}) for several SPEs placed in the 2004–2005 winter. All events begin on 18 December 2004.

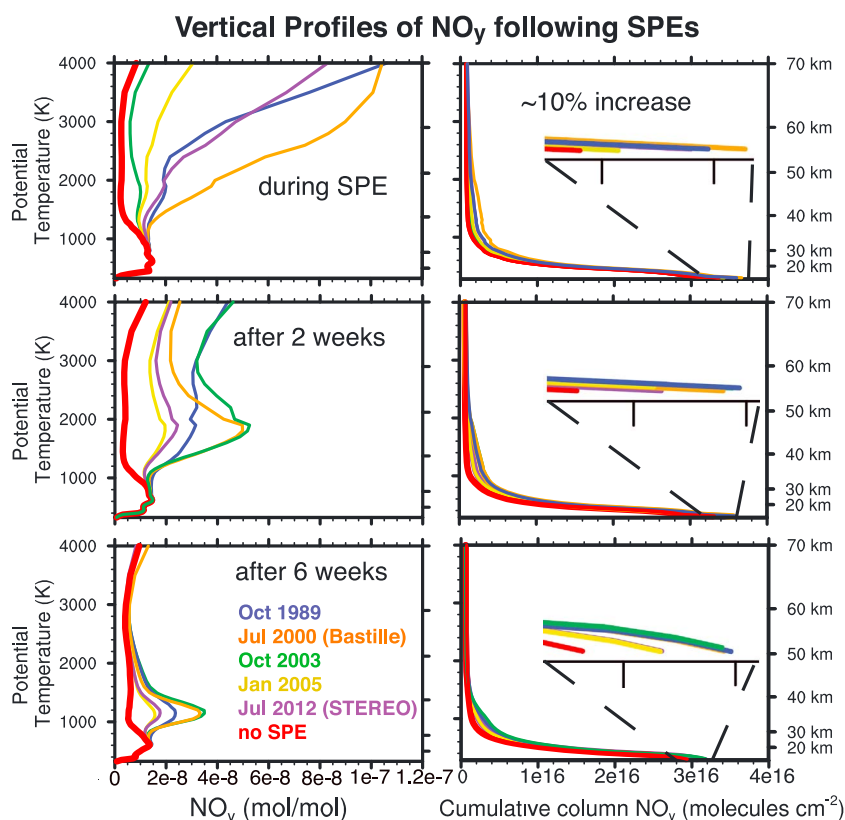


Figure 5. Profiles of Arctic vortex-averaged NO_y (mol/mol) and cumulative column NO_y (molecules cm^{-2}) for several high fluence SPEs placed in the winter of 2004–2005. (top row) During the SPEs. (middle row) Two weeks following SPEs. (bottom row) Six weeks following SPEs.

of January. Nitric acid from PSCs reenters the gas phase below 15–20 km in these WACCM simulations, suggesting that gravitational settling of NAT PSCs and subsequent renitrication near or within the troposphere (generally extending in the Arctic from the surface to ~8–12 km) is unlikely. These results are consistent with renitrication observed during the Polar Aura Validation Experiment of January/February 2005 [Dibb *et al.*, 2006] as well as the previous WACCM analysis of SPEs and NAT PSCs in Duderstadt *et al.* [2014].

Most of the SPE enhancement of NO_y occurs above 30 km. Figure 5 illustrates how relatively large increases in NO_y mixing ratios in the upper atmosphere can lead to minimal enhancement at the surface, a result of the exponential decrease in atmospheric density with altitude. Figure 5 (left column) depicts NO_y enhancements (mol/mol) during, 2 weeks after, and 6 weeks after each event, showing the descent of SPE-enhanced NO_y with time. The corresponding values of cumulative column density (molecules cm^{-2}) demonstrate the minimal impact of SPEs on NO_y in the troposphere, with enhancements at the surface remaining under 15%, consistent with Figure 4.

In the search for SPE-enhanced NO_y discernible above background variability, we increase the fluence of the 19–27 October 1989 SPEs (the largest SPEs with respect to NO_y production) by 10 times, 50 times, and 100 times. Zeller and Dreschhoff [1995] and McCracken *et al.* [2001] attribute a threefold peak in nitrate within the GISP2-H ice core to the 1859 Carrington event (~300 ng/g compared to the summertime maximum of 100 ng/g, with seasonal variations ranging from 20 to 100 ng/g). The largest spikes in the GISP2-H core are 5 times larger than the seasonal cycle amplitude [Zeller and Dreschhoff, 1995; Schrijver *et al.*, 2012]. Smart *et al.* [2014] attribute a GISP2-H nitrate spike 2.5 times background (150 ng/g compared with summertime 60 ng/g) to the 23 February 1956 GLE. Smart *et al.* [2014] also identify a fivefold nitrate ion spike in the BU Summit ice core at a depth corresponding to the 25 July 1946 GLE (~500 ppb in comparison to adjacent ~100 ppb levels) and a threefold spike at a depth dated as corresponding to the 19 November 1949 GLE (~300 ppb in comparison to the summertime peak of ~100 ppb). Based on these examples, our target is to identify SPE fluences that result in twofold to fivefold SPE enhancements of NO_y .

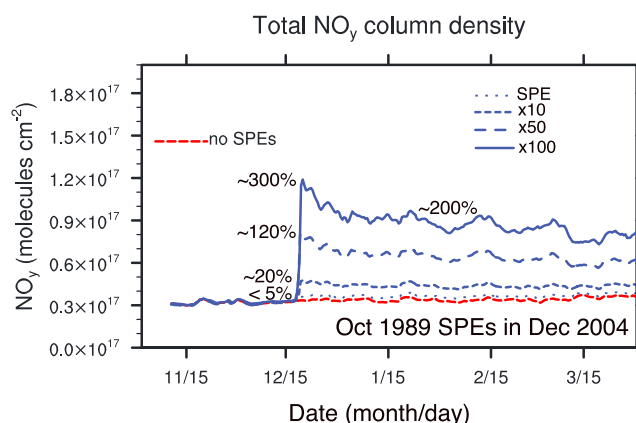


Figure 6. Arctic vortex-averaged total NO_y column density (molecules cm^{-2}) throughout the 2004–2005 winter, comparing simulations without SPEs (red) and scaled simulations with SPEs (blue) based on the October 1989 events.

Total column densities for the extreme SPEs based on the October 1989 events are presented in Figure 6. An event 50 times the October 1989 events would produce total column NO_y enhancements 2 times background levels. An event 100 times the October 1989 events would result in threefold to fourfold enhancements during and following the SPEs, producing enough odd nitrogen throughout the atmospheric column to match the relative magnitude of nitrate spikes attributed to SPEs in the GISP2-H and BU ice cores by Zeller and Dreschhoff [1995], Kepko et al. [2009], and Smart et al. [2014]. This total column NO_y enhancement lasts for several months, suggesting that any potential association with enhanced nitrate deposition would likely occur over a period of months rather than during a single snowfall as indicated by nitrate spikes observed in ice cores

We also consider the production of NO_y directly in the troposphere and lower stratosphere during the October 1989 events, especially given that several GLEs were recorded in a short time period on 19, 22, and 24 October 1989. Figure 7a shows the vortex-averaged column density of NO_x ($\text{NO} + \text{NO}_2$) integrated throughout the troposphere and lower stratosphere (0–15 km) for the October 1989 SPE scaled by 10 times, 50 times, and 100 times. We focus on NO_x in order to capture the maximal immediate SPE enhancement, noting that the lifetime of NO_x with respect to oxidation to NO_y is days to weeks in the troposphere. We choose 0–15 km to remain below the peak background pool of NO_y as well as below the lowest altitude at which PSCs form during the simulation, recalling that PSCs are capable of removing NO_x from the gas phase. Figure 7b shows vortex-averaged 0–15 km column density of NO_y from November to March. It is clear that even with the full spectrum of energies measured by the GOES EPS and HEPAD instruments along with an extrapolated power law tail, there is not enough NO_x or NO_y produced within 0–15 km to suggest a discernible nitrate enhancement at the surface. (Corresponding tropospheric column density enhancements of NO_x and NO_y from 0 to 10 km are less than 1%.)

Figure 7c and 7d show column NO_y from 0 to 30 km and column gas phase HNO_3 plus NAT on PSCs from 15 to 30 km, addressing the potential for enhanced nitrate deposition through gravitational settling of PSCs forming in WACCM from 15 to 30 km during late December through January. The October 1989 $\times 100$ SPE results in up to a 70% enhancement in 0–30 km NO_y during times when PSCs are present in WACCM (recall that NO_y includes nitrate as NAT on PSCs) and a 70% increase in 15–30 km $\text{HNO}_3 + \text{NAT}$ on PSCs.

The short-lived tropospheric NO_x signal is the most reliable indicator for sharp spikes in nitrate deposition, representing direct production in the troposphere. The enhancements of NO_y at all higher levels in the atmosphere lack the characteristics necessary to produce a sharp spike as they last for months as a result of longer lifetimes and slow vertical descent. This behavior includes the gradual enhancement of 0–15 km NO_y that is most likely the result of downward transport from the stratosphere.

4.3. WACCM Simulations of Hard Spectrum Events

From the analysis in the previous section, we conclude that events with the highest F_{30} fluence over the past several solar cycles do not produce NO_y enhancements of similar magnitude to the peaks observed in ice cores, either throughout the atmospheric column or directly in the lower atmosphere. However, the largest

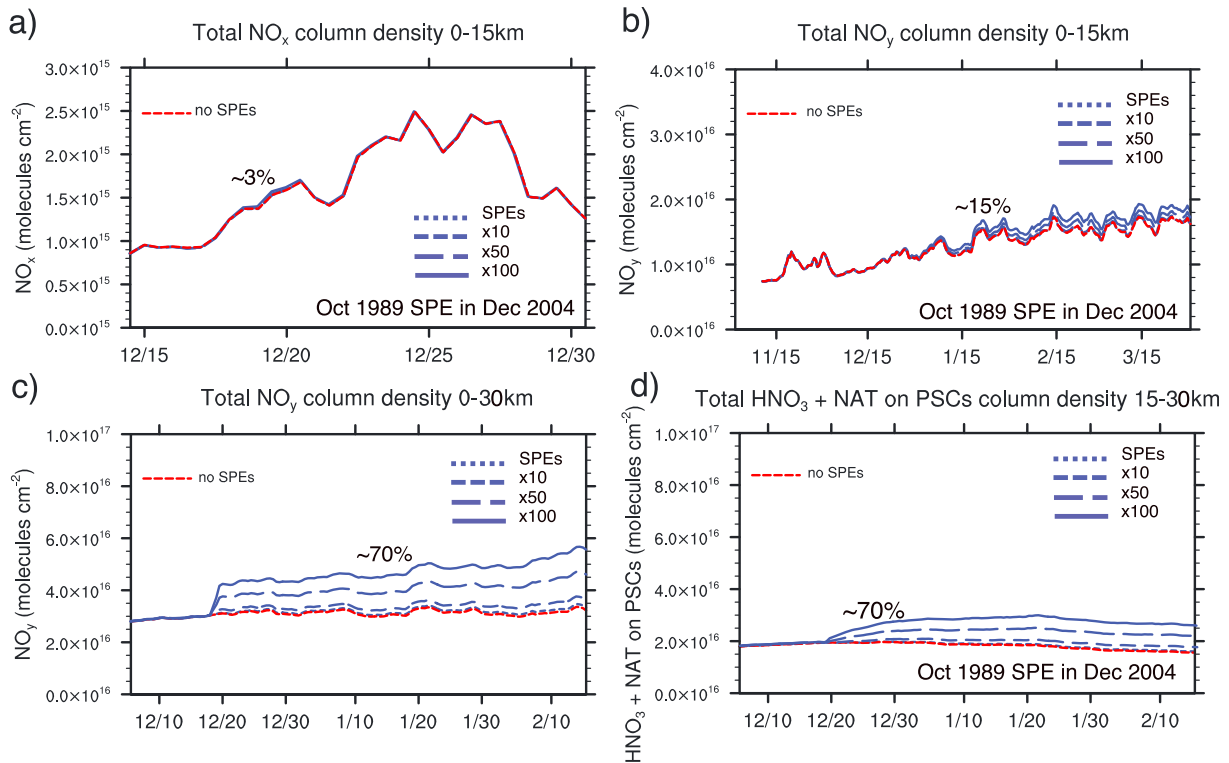


Figure 7. Arctic vortex-averaged column density (molecules cm⁻²) of (a) NO_x from 0 to 15 km, (b) NO_y from 0 to 15 km, (c) NO_y from 0 to 30 km, and (d) HNO₃ from 15 to 30 km. Hypothetical events based on the October 1989 SPEs.

fluence events do not necessarily correspond with the events exhibiting highest energies (hardest spectra). Higher-energy protons are required to produce spallation products that can penetrate and ionize the lower atmosphere. In this section, we investigate the potential for direct production of NO_x at altitudes below the natural background pool of NO_y. This enhanced NO_x would then need to be converted to HNO₃ and deposited to the surface, either directly through wet deposition in the troposphere or via tropopause folds or gravitational settling of polar stratospheric cloud particles in the lower stratosphere.

We base our simulations of hard events on the 29 September 1989 and 20 January 2005 SPEs, with 20 January 2005 having the second largest GLE and 29 September 1989 having the fourth largest GLE since the 1940s [Wang, 2009]. We uniformly scale up the proton fluxes of the 29 September 1989 and 20 January 2005 events until we reach sufficient NO_y enhancements to potentially be observed at the surface. We assume an isotropic proton flux for ion pair calculations but acknowledge observational evidence of a short-lived anisotropy for the 20 January 2005 event [Bazilevskaya et al., 2008].

Figure 8 shows the extrapolated spectral fits to GOES proton data during the prompt components of the 29 September 1989 and 20 February 2005 SPEs. The power law spectral index for the tail of the 20 January 2005 SPE (−2.4 to −2.6) is slightly larger than the *Mewaldt et al.* [2012] fluence spectral index of −2.14 (±0.06). *Mewaldt et al.* [2012] suggest that spectral indices of at least −2.8 are necessary to produce GLEs. *Tylka and Dietrich* [2009] argue that the spectrum steepens for energies greater than ~500 MeV (rigidities greater than ~1 GV), with average integral spectral slopes in rigidity (momentum/charge) of −5 to −7. The extrapolations used in this study consequently represent overestimates of spectral flux, consistent with efforts to determine the minimal characteristics of an SPE necessary to enhance nitrate at the surface.

A more robust estimate of ionization rates and NO_y production requires the addition of galactic cosmic rays (GCRs). While including these particles will increase background ion pair production in the lower atmosphere, it tends to decrease ionization during SPEs as a result of perturbations to the heliospheric magnetic field and/or interplanetary shock; this is known as a Forbush decrease. *Usoskin et al.* [2009], for example, include GCRs in model calculations for the 20 January 2005 SPE, concluding that this event enhances ionization only

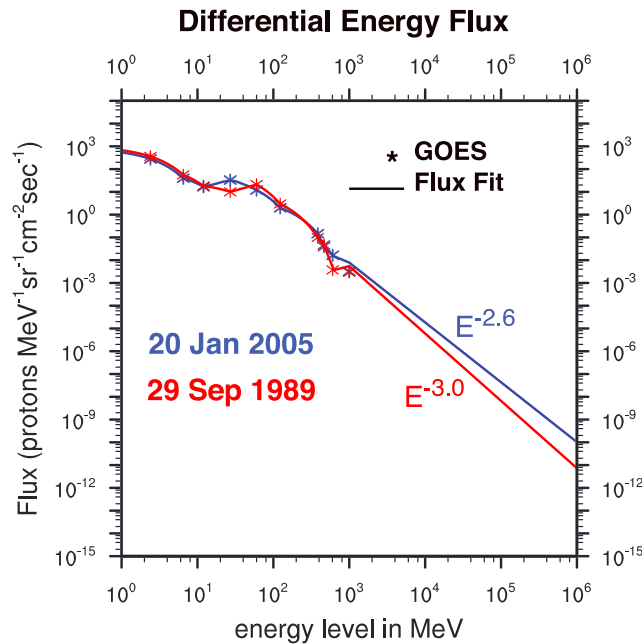


Figure 8. Spectral fits for 20 January 2005 and 29 September 1989 SPEs based on GOES EPS and HEPAD observations. Spectral indices are -3.0 for the 29 September 1989 SPE and -2.6 for the 20 January 2005 SPE.

event (Figure 9a). NO_y SPE enhancements continue to increase throughout the winter, consistent with descending NO_y from the upper stratosphere, peaking at 40% for the 20 January 2005 $\times 1000$ SPE (Figure 9b). Corresponding peak enhancements for the 29 September 1989 SPE simulations are $\sim 2\%$ NO_x and $\sim 25\%$ NO_y for 0–15 km column densities during the 29 September 1989 10d \times 100 SPE (not shown). None of the hypothetical SPEs are capable of producing enhancements of NO_x greater than 1% directly in the troposphere (~ 0 –10 km).

Figures 9c and 9d suggest the potential for twofold enhancements of nitrate deposition from the region impacted by nitrate deposition through PSCs, with WACCM calculations for the hypothetical 20 January 2005 $\times 1000$ and 20 January 2005 10d \times 100 SPEs producing $\sim 130\%$ enhancements of 0–30 km column NO_y and $\sim 150\%$ enhancements of 15–30 km column $\text{HNO}_3 + \text{NAT}$ on PSCs from December to January. Corresponding enhancements for the 29 September 1989 simulations are 60% for 0–30 km column NO_y and 70% for 15–30 km column $\text{HNO}_3 + \text{NAT}$ on PSCs, respectively (not shown).

Figure 10 presents the NO_y total column density for these simulations. An event 1000 times the 29 September 1989 SPE produces a fourfold total column increase in NO_y , which decreases to threefold in the weeks following the event. An event 1000 times the 20 January 2005 SPE produces a fourfold to fivefold increase in NO_y during and following the event. The 10d \times 100 scenario, which has an equivalent total fluence to the $\times 1000$ event, produces a threefold to fourfold increase.

4.4. WACCM Simulations of Carrington-Like Events

Estimates for the September 1859 Carrington event are often used to represent worst-case scenarios in preparation for future solar storms. Most of these studies are based on the $1.9 \times 10^{10} \text{ cm}^{-2} F_{30}$ proton fluence derived from GISP2-H nitrate measurements [McCracken *et al.*, 2001] and aim to quantify the impacts of worst-case scenarios on atmospheric chemistry, temperature, dynamics, and radiation doses. McCracken *et al.* [2001] derive this estimate from a nitrate anomaly in the GISP2-H core dated to the year 1859 by Zeller and Dreschhoff [1995]. Following the Wolff *et al.* [2012] study comparing the GISP2-H core to more robustly dated high-resolution cores, this nitrate spike is likely associated with a 1863 biomass burning event observed in other Greenland cores. This estimate consequently no longer has any observational basis. Nonetheless, we conduct WACCM simulations to determine if this magnitude of F_{30} fluence could produce threefold enhancements in NO_y similar in relative magnitude to the GISP2-H nitrate ion spike. While this

in the middle and upper polar atmosphere, with a Forbush decrease suppressing ionization at lower altitudes and latitudes. Because we want to determine the minimal solar proton flux necessary to achieve maximal enhancement of NO_y , we have chosen not to include GCRs in these hard spectra simulations.

We perform a series of case studies increasing the fluxes of the 29 September 1989 and 20 January 2005 SPEs as well as extending the duration of the events. Figure 9 presents WACCM calculations of vortex-averaged 0–15 km NO_x and NO_y , 0–30 km NO_y , and 15–30 km HNO_3 for the 20 January 2005 SPE, with an analysis rationale similar to Figure 7 described in section 4.2. The 20 January 2005 10d \times 100 SPE results in the greatest direct enhancement of NO_x in the troposphere and lower stratosphere, peaking at 10% during the

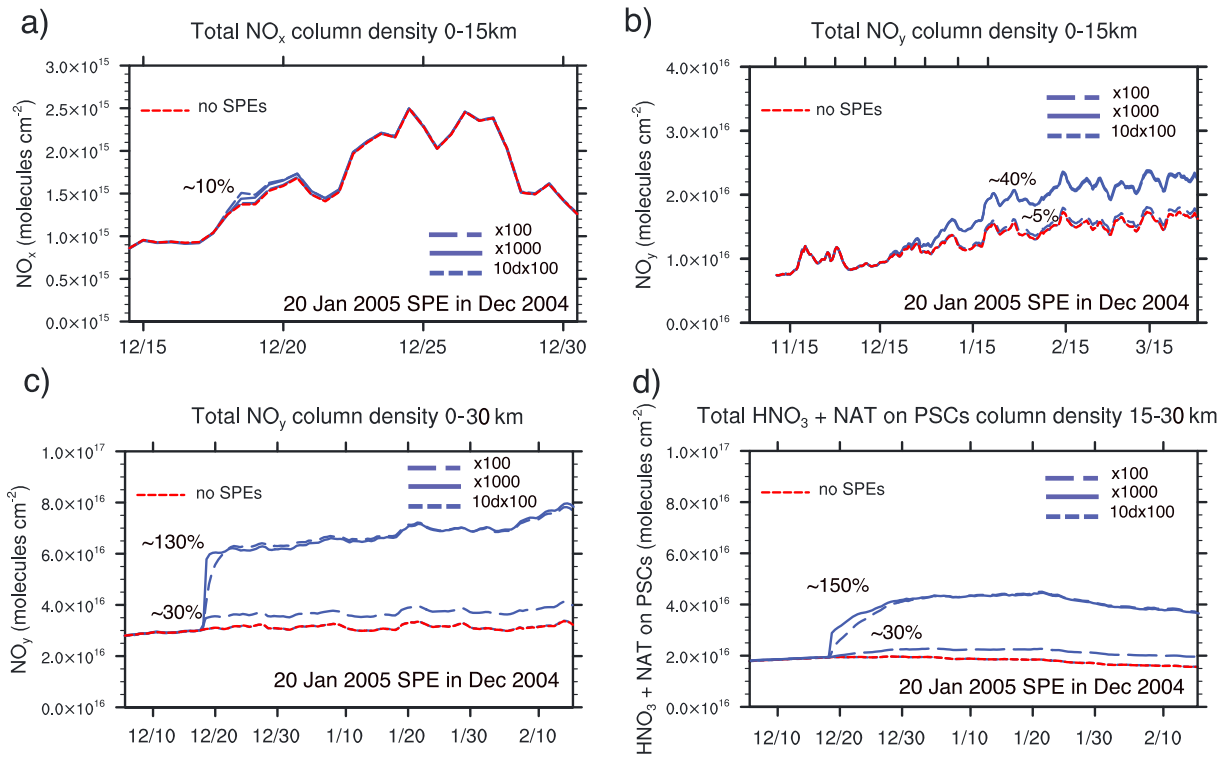


Figure 9. Arctic vortex-averaged column density (molecules cm⁻²) of (a) NO_x from 0 to 15 km, (b) NO_y from 0 to 15 km, (c) NO_y from 0 to 30 km, and (d) HNO₃ from 15 to 30 km. Hypothetical events based on the 20 January 2005 SPEs.

approach is somewhat circular, it provides a simple means to determine if the underlying assumption of the $1.9 \times 10^{10} \text{ cm}^{-2} F_{30}$ used in prior studies to estimate the fluence for a worst-case scenario is well founded.

There are many examples of prior Carrington-like event simulations. *Townsend et al.* [2003, 2006] use a Weibull fit for the 4 August 1972, August-September-October 1989, and 23 March 1991 SPEs scaled to the $1.9 \times 10^{10} \text{ cm}^{-2} F_{30}$ fluence to represent the Carrington event. *Rodger et al.* [2008] consider Weibull fits for the 4 August 1972 and 23 March 1991 SPEs scaled to the $1.9 \times 10^{10} \text{ cm}^{-2} F_{30}$ fluence as well, arguing that soft energy spectra are most representative of the Carrington event due to the lack of signature for the event in ¹⁰Be ice core data. *Miroshnichenko and Nymmik* [2014] take the average spectrum from five SPEs with the largest F_{30} fluence from solar cycle 23 (8 November 2000, 24 September 2001, 4 November 2001, 28 October 2003, and 17 January 2005) to obtain a differential spectrum representative of the Carrington event, scaling it once again to

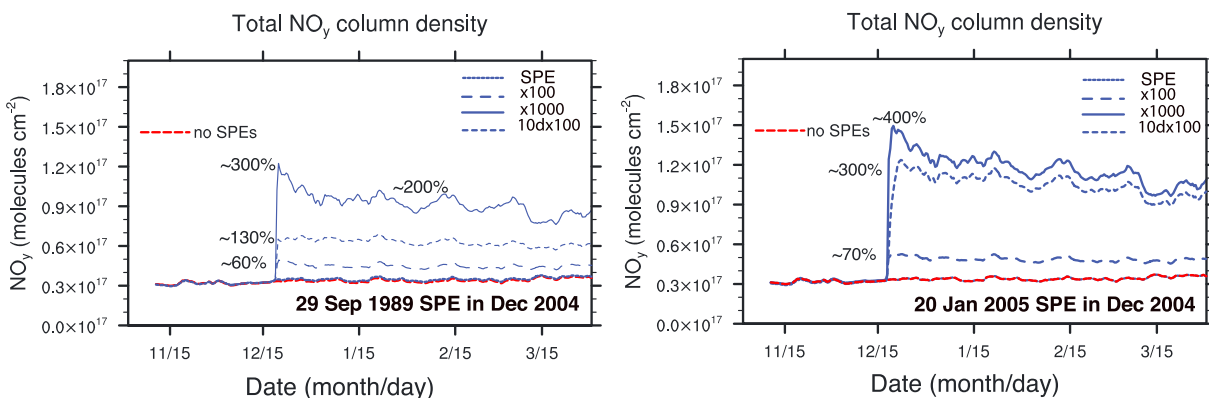


Figure 10. Arctic vortex-averaged total NO_y column density (molecules cm⁻²) throughout the 2004–2005 winter, comparing simulations without SPEs (red) and simulations with SPEs (blue) scaled from the 29 September 1989 and 20 January 2005 SPEs.

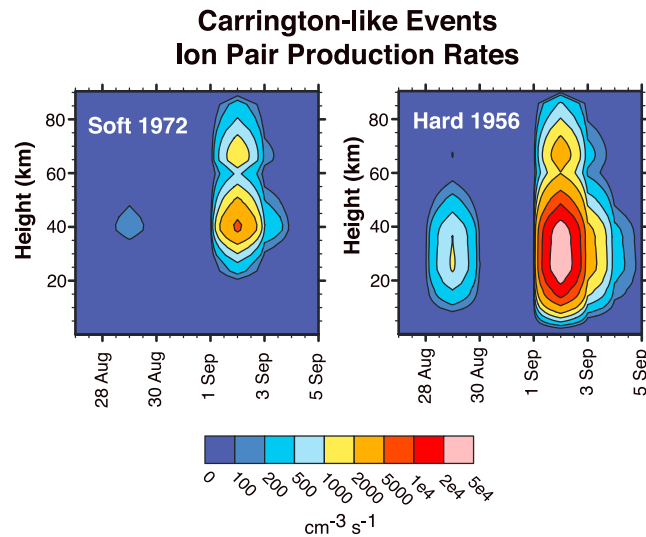


Figure 11. Ionization profiles for Carrington-like events as a function of time. (left) Soft Carrington-like event modeled after August 1972 SPE. (right) Hard Carrington-like event modeled after February 1956 SPE.

the $1.9 \times 10^{10} \text{ cm}^{-2} F_{30}$ fluence. *Calisto et al.* [2012] similarly scale the 4 August 1972 event to the F_{30} fluence of $1.9 \times 10^{10} \text{ cm}^{-2} F_{30}$, using the time profile reconstructed by *Smart et al.* [2006a] and the three-dimensional Chemistry Climate Model SOCOL v2.0 to study atmospheric impacts. *Calisto et al.* [2013] improve on this effort to study a worst-case scenario by basing a Carrington-like event on the hard 23 February 1956 SPE and adding GCRs. The authors calculate ionization rates for high-energy protons using the CRAC:CRIL particle transport model. *Calisto et al.* [2013] find statistically significant SPE influences not only on upper atmospheric HO_x , NO_x , and O_3 but also on zonal winds and surface tem-

peratures. The authors calculate up to a 40% increase in HNO_3 deposition during the month following the Carrington-like 1956 simulations in regions over the South Pole but note a less pronounced effect in the Arctic.

We conduct similar Carrington-like simulations based on the August 1972 and February 1956 proton fluxes, using the ionization profiles of *Usoskin et al.* [2010] and *Calisto et al.* [2013]. Figure 11 shows the ionization rates used as input for WACCM. These simulations include ion pair production rates from GCRs as described in *Calisto et al.* [2013]. Figure 12 presents a comparison of the peak ionization rate (IR) profiles used in WACCM for the Carrington-like soft (1972) and hard (1956) simulations along with the October 1989, January 2012, and March 2012 rates calculated using 1–300 MeV protons [*Jackman et al.*, 2014]. This figure also shows the August 1972 ionization rates from *Rodger et al.* [2008] and *Calisto et al.* [2012] and annual mean GCR rates by *Mertens et al.* [2013]. It is clear that the hard Carrington-like event based on the 23 February 1956 SPE results in the greatest levels of total ionization, with significantly more ionization in the lower atmosphere than background GCRs (1 to 2 orders of magnitude).

As a comparison with the previous model runs, Figure 13 presents the vortex-averaged NO_y total column density for both the Carrington-like soft (Soft Car 1972) and Carrington-like hard (Hard Car 1956) events, where the reconstructed 1 September Carrington peak is placed on 13 December 2004. The Soft Car 1972 event results in a 5% total column increase in NO_y , while the Hard Car 1956 event produces a 20% total column increase in NO_y . This result is not surprising considering the greater atmospheric ionization rates of the Hard Car 1956 event at lower altitudes, where atmospheric densities are much higher. These Carrington-like events do not produce enough odd nitrogen within the atmospheric column to account for the threefold spike in nitrate ions attributed to the Carrington event in the GISP2-H ice core.

As a result of the Hard Car 56 simulated event, WACCM calculates a ~3% maximal increase in NO_x vortex-averaged column densities for 0–10 km and ~15% for 0–15 km during the event, with <10% enhancements of 0–15 km column density throughout winter. Enough NO_x is produced in the atmosphere to result in a 1% reduction of O_3 at 20 km and 10–20% reduction of O_3 from 30 km to 40 km in the weeks following the event, consistent with the results of *Calisto et al.* [2013]. However, even with rapid oxidation of NO_x to HNO_3 and strong stratosphere-tropospheric exchange, this enhancement would not be enough to account for the nitrate peaks observed in ice cores. The potential impact of gravitational settling of nitrate on PSCs is also minimal, with enhancements of 10% for 0–30 km column densities of NO_y (including NAT on PSCs) and 15% for 15–30 km column densities of HNO_3 + NAT on PSCs.

It is unlikely that a $1.9 \times 10^{10} \text{ cm}^{-2} F_{30}$ fluence could produce a threefold enhancement of NO_y in the lower atmosphere similar to the nitrate spike in the GISP2-H ice core attributed to the Carrington event.

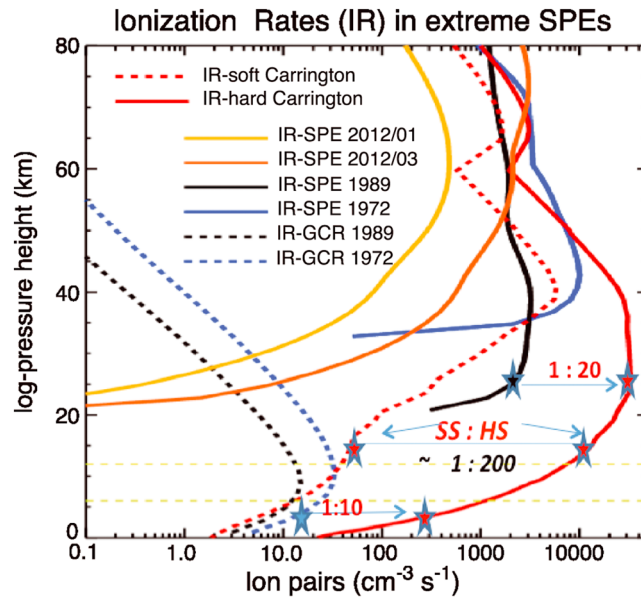


Figure 12. Comparison of ionization profiles of the “IR-soft” (August 1972-based) and “IR-hard” (February 1956-based) Carrington-like SPEs used in WACCM simulations with SPE profiles from August 1972, October 1989, January 2012, and March 2012 and GCR rates (annual mean) from 1972 and 1989 [Usoskin et al., 2010, 2012; Mertens et al., 2013; Jackman et al., 2014].

However, as we discuss in the next section, satellite and cosmogenic radionuclide studies show that a F_{30} fluence of this magnitude is within the occurrence probability of a millennium timescale, supporting prior studies involving worst-case scenarios based on this fluence.

5. Discussion

Our WACCM experiments indicate that events 50 times and 100 times larger than the 1989 October series of events would be required to increase NO_y by twofold and threefold to fourfold, respectively, throughout the atmosphere. An event 3 orders of magnitude larger than the hard spectrum 20 January 2005 and 29 September 1989 events would be necessary to produce a short-lived twofold to threefold increase of NO_x at altitudes below the natural pool of stratospheric NO_y . These enhancements are similar in relative magnitude to the twofold to fivefold nitrate spikes measured in Greenland ice cores and attributed to SPEs.

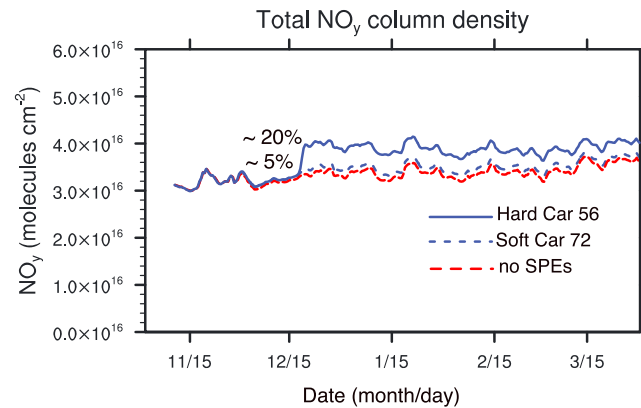


Figure 13. Arctic vortex-averaged total NO_y column density (molecules cm^{-2}) throughout the 2004–2005 winter, comparing simulations without SPEs (red) and Carrington-like simulations based on the August 1972 SPE (Soft Car 72) and February 1956 SPE (Hard Car 56).

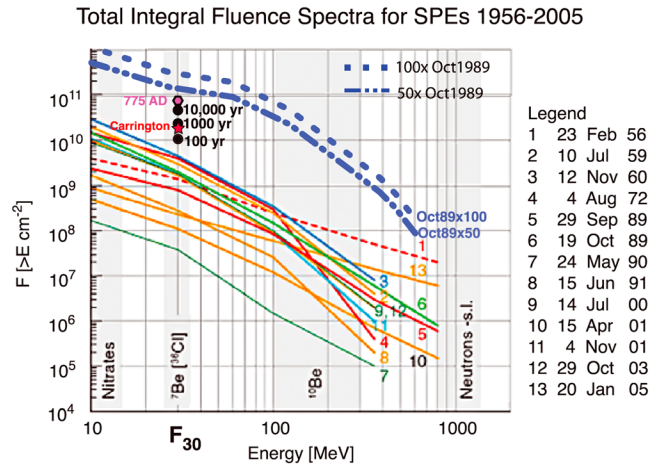


Figure 14. Total integral fluence spectra for SPEs from 1956 to 2005 adapted from Webber et al. [2007] and Beer et al. [2012]. Shaded regions indicate peak response energies for nitrates, ⁷Be, ¹⁰Be, and neutron monitors. Hypothetical events for 50X and 100X the October 1989 SPEs used in WACCM simulations are overlaid on the plot. Black circles indicate probability of occurrence from Usoskin and Kovaltsov [2012]. Red star identifies the estimated Carrington event fluence from McCracken et al. [2001] based on GISP2-H nitrate. Purple hexagon shows fluence for potential SPE in 774–775 [Cliver et al., 2014].

from Webber et al. [2007] and Beer et al. [2012] for events from 1956 to 2005. This figure illustrates how extreme an SPE would have to be in order to produce 2 to 4 times the amount of total column NO_y. These hypothetical events are well above the integral fluences calculated for measured events. The red star represents the F₃₀ estimate of 1.9 × 10¹⁰ cm⁻² for the Carrington event from McCracken et al. [2001] used in section 4.4, well below the magnitude of the October 1989 ×50 SPE needed to produce a doubling of NO_y throughout the atmosphere.

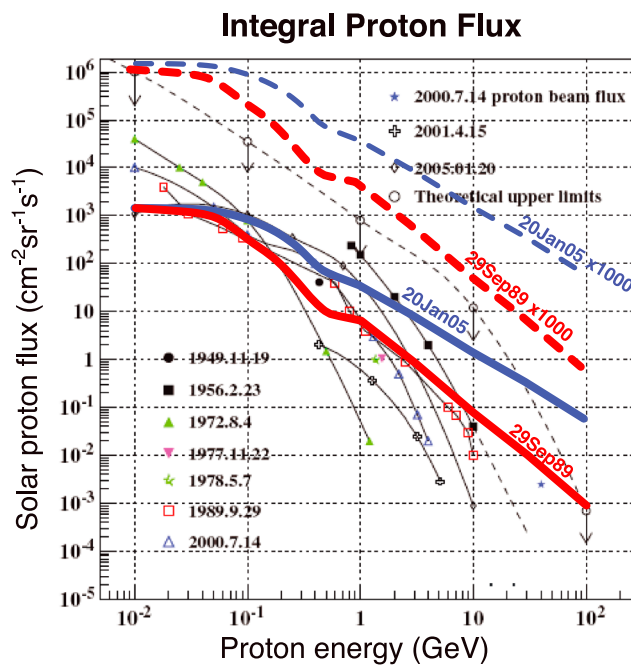


Figure 15. Integral flux of solar proton events during peak flux. Adapted from Wang [2009]. Hard spectra events used in WACCM are overlaid in red (29Sep89 and 29Sep89x1000) and blue (20Jan05 and 20Jan05x1000).

By placing these hypothetical SPEs in the context of probability distributions and upper limit estimates, we can assess the frequency that such events might occur and consequently the potential usefulness of nitrate as a proxy for individual SPEs. It is important to keep in mind that we are considering only the contribution of odd nitrogen enhancement in the atmosphere. There are many additional arguments challenging the use of nitrate as a proxy for SPEs that are beyond the scope of this paper, including stratospheric residence time, vertical and latitudinal mixing, alternative sources of nitrate, sequestration of nitrate within buried ice, and postdepositional processing.

Figure 14 presents total integrated fluences for the hypothetical October 1989 ×50 and October 1989 ×100 SPEs overlaid onto fluence summaries

Several studies calculate the probability of occurrence of SPEs using satellite measurements, GLEs, and terrestrial and lunar cosmogenic isotopes. The slope of the cumulative probability of occurrence becomes steeper at larger fluences, resulting in a broken power law distribution [Smart et al., 2006b]. Usoskin and Kovaltsov [2012] analyze cosmogenic isotopes ¹⁴C and ¹⁰Be in terrestrial archives to determine the probability of occurrence of SPEs with respect to F₃₀. They calculate conservative F₃₀ limits of 1 × 10¹⁰ cm⁻², 2–3 × 10¹⁰ cm⁻², and 5 × 10¹⁰ cm⁻² occurring on the order of 100 years, 1000 years, and 10,000 years, respectively. These levels are plotted as black circles in Figure 14 and are below the F₃₀ fluences of our hypothetical scenarios. Kovaltsov and Usoskin [2014] arrive at a cumulative occurrence probability distribution function based on 60 years of direct

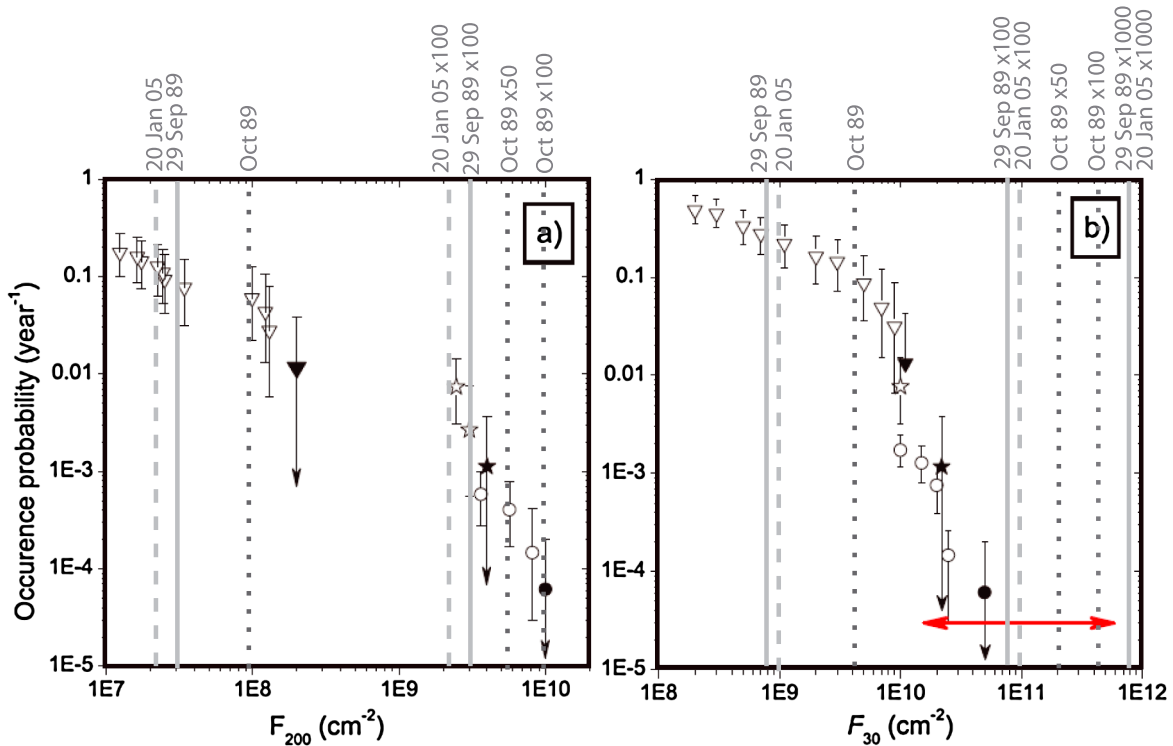


Figure 16. Probability of occurrence for SPEs having fluence integrated above (a) 200 MeV (F_{200}) and (b) 30 MeV (F_{30}). Adapted from *Usoskin and Kovaltsov* [2012] and *Kovaltsov et al.* [2014]. The red arrow indicates the range of uncertainty due to unknown spectral shape as in *Kovaltsov et al.* [2014]. Vertical gray lines represent the hypothetically extreme SPEs used in our WACCM simulations.

fluence measurements, terrestrial ^{14}C and ^{10}Be records during the Holocene, and cosmogenic radionuclides in lunar rocks. The authors confirm that SPEs with F_{30} fluences greater than 10^{11} protons $\text{cm}^{-2}\text{yr}^{-1}$ are not expected on a millennium timescale, and no events greater than 5×10^{11} protons yr^{-1} are expected to have occurred during the past 10,000 years. The probabilities at higher fluences of *Usoskin and Kovaltsov* [2012] and *Kovaltsov et al.* [2014] rely on the identification of ^{14}C and ^{10}Be peaks in ice core archives, with significant uncertainty regarding the contributions from Earth system effects and local meteorology. Nonetheless, cosmogenic radionuclides remain the current best estimate. *Miroshnichenko and Nymmik* [2014] evaluate the fluence of events measured by satellites during recent solar cycles combined with historical sunspot data to conclude that an event with F_{30} greater than $6 \times 10^{10} \text{ cm}^{-2}$ is only likely to occur once every 2.6×10^5 years.

The October 1989 series of events (19–27 October) has a total F_{30} fluence of $\sim 4 \times 10^9 \text{ cm}^{-2}$ (Table 1). An event 50 times the October 1989 events, necessary to double polar total column odd nitrogen, would consequently have a F_{30} fluence of $\sim 2 \times 10^{11} \text{ cm}^{-2}$, which is well beyond the Holocene probability estimates from *Usoskin and Kovaltsov* [2012] and *Miroshnichenko and Nymmik* [2014]. An event 100 times the October 1989 event would have a F_{30} fluence of $\sim 4 \times 10^{11} \text{ cm}^{-2}$, approaching the Holocene upper limit indicated by *Kovaltsov and Usoskin* [2014]. *Usoskin et al.* [2013] estimate that an extremely large SPE capable of producing the spikes observed in 774–775 A.D. cosmogenic radionuclide archives [*Miyake et al.*, 2012, 2015] would be 25–50 times stronger than the February 1956 SPE which had a F_{30} fluence of $\sim 8 \times 10^{10} \text{ cm}^{-2}$ as estimated by *Cliver et al.* [2014] (purple hexagon in Figure 14). *Thomas et al.* [2013] assert that a smaller event, an order of magnitude greater than the October 1989 SPEs, would explain the 774–775 spike. Both estimated fluences are less than the amount required by WACCM to produce a doubling of odd nitrogen in the atmospheric column, implying that signatures of a 774–775 A.D. event are not expected in ice core nitrate records. In addition, *Cliver et al.* [2014] question an event of this size, suggesting that it would require an active region on the Sun (sunspot size) 2.5 times larger than has been recorded, further challenging the likelihood of the extreme hypothetical October 1989 $\times 50$ and October 1989 $\times 100$ events. This argument is consistent with the *Aulanier et al.* [2013] estimate for an upper limit to solar flare energies equivalent to 6 times the 4 November 2003 flare. The implications of this analysis are that SPE-

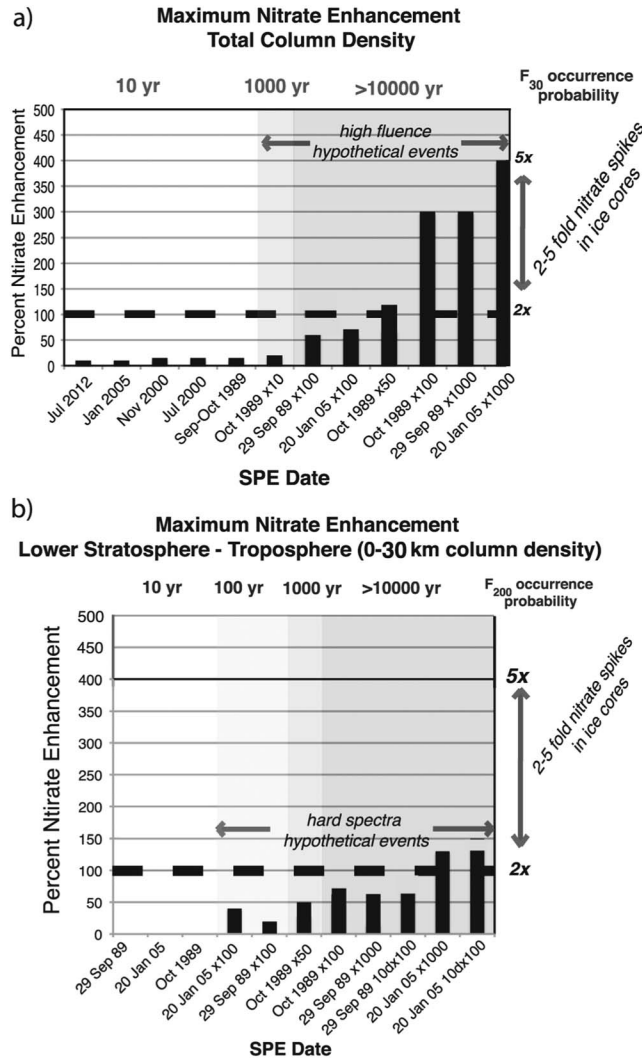


Figure 17. Nitrate enhancement resulting from SPEs measured by satellite and hypothetical events. Probability of occurrence from Kovaltsov *et al.* [2014] indicated on top horizontal axis and as shaded regions. (a) Peak total column density of NO_y for high fluence events with F₃₀ probability of occurrence. (b) Peak 0–30 km NO_y column density for hard spectrum events with F₂₀₀ probability of occurrence.

enhanced odd nitrogen in the mesosphere and upper stratosphere is insufficient to produce discernible nitrate spikes at the surface for any meaningful analysis of individual SPE statistics during the Holocene.

The black dashed line in Figure 15 represents a theoretical upper limit spectrum [Miroshnichenko and Nymmik, 2014], calculated by averaging the time-of-maximum intensity spectra of the five largest F₃₀ SPEs of solar cycle 23 and scaling to a fluence an order of magnitude larger than the 1.9 × 10¹⁰ cm⁻² McCracken *et al.* [2001] estimation for the Carrington event. We note that although the fluence of this upper limit spectrum is determined using a nitrate source that has been challenged by Wolff *et al.* [2012], the shape of the spectrum is determined from direct measurements and is remarkably consistent with the shape of our hypothetical events from 30 MeV to 10 GeV. The peak fluxes for the hypothetical 29 September 1989 × 1000 and 20 January 2005 × 1000 events exceed all other events during the satellite era as well as the Miroshnichenko and Nymmik [2014] upper limit spectrum.

Figure 16 shows probability of occurrence distribution functions for F₃₀ and F₂₀₀ from Usoskin and Kovaltsov [2012] and Kovaltsov *et al.* [2014] based on satellite data and cosmogenic radionuclides. The vertical gray lines indicate the fluences of the hypothetical events in the WACCM simulations. The 29 September 1989 × 100 and 20 January 2005 × 100 hypothetical SPEs, producing only a 15–40% enhancement of NO_x in the lower stratosphere, occur on the order of once every 100–1000 years and would likely not be distinguishable from seasonal variability. Usoskin and Kovaltsov [2012] use ¹⁰Be data to identify 19 candidates throughout the Holocene having 10–30 times the fluence of the 23 February 1956 SPE. These events, with F₂₀₀ on the order

Next, we assess the probability of hard spectrum SPEs producing odd nitrogen directly in the lower atmosphere by analyzing the total integral fluence above 200 MeV (defined to be F₂₀₀) as well as the instantaneous flux of protons during peak intensity of our hypothetical hard spectrum SPEs. Figure 15 is adapted from Wang [2009] and summarizes instantaneous solar proton integral flux as a function of energy during times of peak intensity for several measured events. Overlaid onto this figure are the integral flux spectra during the prompt component of the 29 September 1989 SPE and 20 January 2005 SPE used in our WACCM simulations along with the corresponding 1000-fold hypothetical events. Wang [2009] calculate spectra using neutron monitor data, ionization chamber data, and GOES and Meteor direct satellite measurements. Our simulations use power law fits from 30 min average GOES data, resulting in larger flux at relativistic energies as discussed in section 2.

The black dashed line in Figure 15 represents a theoretical upper limit spectrum [Miroshnichenko and Nymmik, 2014], calculated by averaging the time-of-maximum intensity spectra of the five largest F₃₀ SPEs of solar cycle 23 and scaling to a fluence an order of magnitude larger than the 1.9 × 10¹⁰ cm⁻² McCracken *et al.* [2001] estimation for the Carrington event. We note that although the fluence of this upper limit spectrum is determined using a nitrate source that has been challenged by Wolff *et al.* [2012], the shape of the spectrum is determined from direct measurements and is remarkably consistent with the shape of our hypothetical events from 30 MeV to 10 GeV. The peak fluxes for the hypothetical 29 September 1989 × 1000 and 20 January 2005 × 1000 events exceed all other events during the satellite era as well as the Miroshnichenko and Nymmik [2014] upper limit spectrum.

of 10^9 cm^{-2} and F_{30} on the order of 10^{10} cm^{-2} , would be well below the limits needed to produce a doubling of odd nitrogen in the lower atmosphere or throughout the atmospheric column. Within the *Usoskin and Kovaltsov* [2012] range of uncertainty in spectral shape (indicated by the red arrow in Figure 16b), the October 1989 $\times 100$, 29 September 1989 $\times 1000$, and 20 January 2005 $\times 1000$ events would take place at most, on average, once every 12,500 years (occurrence probability of 8×10^{-5}), implying that SPEs of this size are unlikely to have occurred during the Holocene.

6. Conclusion

This study sets out to quantify the fluence, duration, and energy spectra of solar proton events necessary to produce enhancements of NO_y equivalent to nitrate spikes in the GISP2-H and BU ice cores that have been attributed to SPEs. We place these extreme events in the context of probability of occurrence distributions constructed using measured events, sunspot records, and cosmogenic radionuclide data. We conclude that events necessary to produce nitrate spikes in ice cores are beyond the probability of occurrence during the Holocene, confirming that nitrate ion spikes cannot be used as one to one, or even statistically representative, proxies for solar proton events.

Figure 17 summarizes our results. Figure 17a shows the total column NO_y enhancement from large SPEs during the satellite era as well as high fluence hypothetical events, with the dashed horizontal line indicating a twofold increase. The top horizontal axis and shading show occurrence probabilities based on integral proton fluence greater than 30 MeV. An event 50 times the October 1989 series of events would be required to enhance NO_y in the atmospheric column by a factor of 2. It would take an event 100 times the October 1989 series of events to produce a threefold to fourfold increase in NO_y , similar in relative magnitude to the nitrate ion spikes in Greenland ice cores that some have attributed to SPEs. Probability distributions suggest that events of this size, with total integral fluence over 30 MeV of $\sim 2 \times 10^{11} \text{ cm}^{-2}$ (October 1989 $\times 50$) and $\sim 4 \times 10^{11} \text{ cm}^{-2}$ (October 1989 $\times 100$), are unlikely to have occurred during the Holocene ($\sim 10,000$ years). Even rare millennial timescale events such as the event suggested by ^{14}C and ^{10}Be in 774–775 A.D. would be indiscernible in the nitrate record.

Figure 17b summarizes results from hard spectra (high-energy) SPEs capable of producing odd nitrogen directly in the lower atmosphere. These graphs show NO_y column enhancements from 0 to 30 km during the time period when PSCs are present, accounting for the SPE enhancement of odd nitrogen in the troposphere and lower stratosphere as well as the potential for enhanced nitrate deposition through uptake by gravitational settling of PSCs. Only the 20 January 2005 $\times 1000$ and 20 January 2005 $10d \times 100$ hypothetical SPEs produce a doubling of 0–30 km NO_y . Enhancements of NO_x and NO_y directly in the troposphere (0–10 km) are less than 1% for all simulations throughout the winter, while enhancements reaching the lower stratosphere (0–15 km) are less than 10% for NO_x and less than 40% for NO_y for the hypothetical SPEs. While events 3 orders of magnitude larger than the hard spectrum 20 January 2005 and 29 September 1989 SPEs show minimal significant increase in tropospheric NO_x , they are capable of producing a 130% and 60% increase of NO_y from 0 to 30 km, respectively. Such hypothetical events, however, are well beyond the limits of the probability of occurrence during the Holocene.

These results show, in the framework of the present knowledge of atmospheric processes, that nitrate spikes in ice cores are not associated with individual SPEs. We conclude that nitrate ions in ice cores cannot, therefore, be used as proxies to study the frequency and strength of historical SPEs. Future efforts to analyze new cores for higher-resolution nitrate spikes in the search for statistics involving SPEs are unwarranted. The results substantiate the argument by *Wolff et al.* [2012] that the 1859 Carrington SPE is not observed in ice cores and support the exclusion of nitrate as a proxy in the study of the frequency of extreme solar events by *Schrijver et al.* [2012]. Nitrate ion spikes observed in ice cores are not suitable proxies for SPEs.

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